

Mukai-Fourier Transform in Derived Categories to Solutions of the Field Equations: Gravitational Waves as Oscillations in the Space-Time Curvature/Spin IV

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Key words: Classical limit conjecture, derived deformed categories, Fourier-Mukai equivalence, Hecke functors, Higgs Bundel, Hamiltonian variety, Langlands correspondence, Mukai-Furier transform.

UDC:511 512 514.7 515.1 517 517.9

2010 AMS Classification: 53D37; 11R39; 14D24; 83C60; 11S15.

Starting of fact that the Mukai-Fourier transform is an equivalence of derive categories(with arbitrary decorations: +, -,b), is feasible construct a Fourier-Mukai equivalence given for

$$D_{Coh}(T^\vee A) \cong D_{Coh}(A^\vee \times \mathbf{H})$$

, where exist a distinguished deformation of the category $D_{Coh}(T^\vee A)$, which is a non-commutative deformation of $T^\vee A$, defined by a natural symplectic form, that is their quatization [1].

Then $T^\vee A$, results a 1-parameter deformation A^b , of the space $A^\vee \times \mathbf{H}$, to an affine bundle over A^\vee , classified by $H^1(A^\vee; \mathcal{O} \otimes \mathbf{H})$. Then the Fourier-Mukai equivalence relative to the projection $T^\vee A$, deforms an equivalence between the deformed categories $D_{Coh} \mathcal{D}_A - mod$, and $D_{Coh}(A)^b$.

Then we use the deformed version of the Mukai-Fourier transform that results on D_A- modules and we characterize to A , as a Picard variety of C , ¹, where C , is a curve. Then a Hecke functor is definid as the integral transform

$$\Phi^1 : D_{Coh}(Pic(C), \mathcal{D}) \rightarrow D_{Coh}(C \times Pic(C), \mathcal{D}),$$

to D -modules on ${}^L Bun$. But using the classical limit conjeture is had the equivalence through of the interpretation of Higgs sheaves, given in the category $D_{Coh}({}^L Higgs_0, \mathcal{O})$, which can be extended to the corresponding Langlands correspondence \mathfrak{c} , of the quantum sheaves given by $\mathfrak{c} = quant_{Bun} \circ \Phi \circ quant_C^{-1}$, where Φ , is the Fourier-Mukai transform that we want. Then we have as integral the integral transforms composition [2] $\mathfrak{c} \circ \Phi^\mu = {}^L \Phi^\mu$, which is solution to the field equation $Isom d\mathbf{h} = 0$, where \mathbf{h} , are the cotangent vector (Higgs fields).

Then by superposing of these states, considering the field corresponding ramifications(connections), we have

$$\mathcal{H} = \mathbf{H}^0(\omega_c) \oplus \mathbf{H}^0(\omega_C^{\otimes 2}) \oplus \dots \oplus \mathbf{H}^0(\omega_C^{\otimes n}),$$

which has their re-interpretation as the curvature energy expressed through the H-states which can be written using the superposing principle to each connection $\omega_C^{\otimes j}$, (with C , a curve) that describes the corresponding dilaton (photon or gauge particle).

Likewise, in a Hamiltonian densities space [3] we have the Figure 1, considering a Hitchin base. In the case of a spinor representation the corresponding H-states can be given as spinor waves (Figure 2) which can be consigned in oscillations in the space-time-curvature/spin, to a microscopic deformation measured [4] in \mathcal{H} .

¹In a physical context (could be taken $\mathbb{M} = Pic(C)$, where \mathbb{M} , is the space-time), this represent a trace of particles in the symplectic geometry that can be characterized in a Hamiltonian manifold.

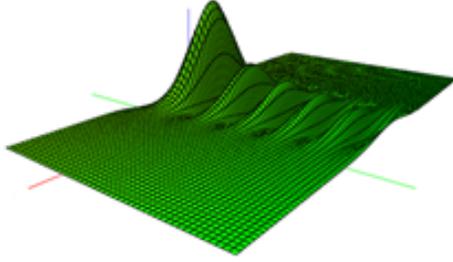


Figure 1

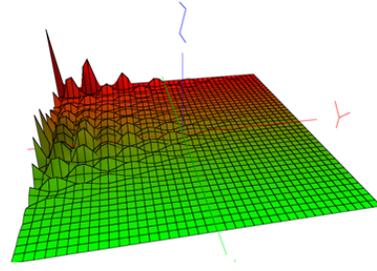


Figure 2

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