

# Dynamics and exact solutions of linear PDEs

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The report presents a new method for constructing exact solutions of the classical linear equations of mathematical physics of parabolic, hyperbolic, elliptic and variable types. The method is a generalization of the theory of finite-dimensional dynamics proposed for evolutionary differential equations [1, 5]. The theory of finite-dimensional dynamics is a natural development of the theory of dynamical systems. Dynamics make it possible to find families that depends on a finite number of parameters among all solutions of PDEs (see [2, 3]).

Consider the following class of second order linear partial differential equations

$$u_{tt} + 2b(x)u_{tx} + c(x)u_{xx} + h(x)u_t + g(x)u_x + f(x) = 0, \quad (1)$$

where  $b, c, h, g, f$  are functions of the class  $C^\infty$ . Such equations are equivalent to the following evolutionary systems

$$\begin{cases} u_t = v, \\ v_t = -2b(x)v_x - c(x)u_{xx} - h(x)v - g(x)u_x - f(y). \end{cases} \quad (2)$$

We call the system of ordinary differential equations of order  $k + 1$

$$\begin{cases} y^{(k+1)} = Y(x, y, z, y', z' \dots, y^{(k)}, z^{(k)}), \\ z^{(k+1)} = Z(x, y, z, y', z' \dots, y^{(k)}, z^{(k)}) \end{cases} \quad (3)$$

a *dynamics* of equation (1) if the vector function

$$(\varphi, \psi) := (z_0, -2b(x)z_1 - c(x)y_2 - h(x)z_0 - g(x)y_1 - f(x))$$

is a generating function of infinitesimal characteristic symmetries of this system [4]. Here  $x, y_0, z_0, y_1, z_1, y_2, z_2$  are canonical coordinates on the space of 2-jets  $J^2(\mathbb{R}^1, \mathbb{R}^2)$ .

**Theorem 1.** *The vector field on  $J^k(\mathbb{R}^1, \mathbb{R}^2)$*

$$S = \varphi \frac{\partial}{\partial y_0} + \psi \frac{\partial}{\partial z_0} + \mathcal{D}(\varphi) \frac{\partial}{\partial y_1} + \mathcal{D}(\psi) \frac{\partial}{\partial z_1} + \dots + \mathcal{D}^k(\varphi) \frac{\partial}{\partial y_k} + \mathcal{D}^k(\psi) \frac{\partial}{\partial z_k} \quad (4)$$

is an infinitesimal characteristic symmetry of system (3) if the following conditions hold:

$$\begin{cases} \mathcal{D}^{k+1}(\varphi) - S(Y) = 0, \\ \mathcal{D}^{k+1}(\psi) - S(Z) = 0. \end{cases} \quad (5)$$

Here

$$\mathcal{D} = \frac{\partial}{\partial x} + y_1 \frac{\partial}{\partial y_0} + z_1 \frac{\partial}{\partial z_0} + \dots + y_k \frac{\partial}{\partial y_{k-1}} + z_k \frac{\partial}{\partial z_{k-1}} + Y \frac{\partial}{\partial y_k} + Z \frac{\partial}{\partial z_k}.$$

Let  $\Gamma^k \subset J^2(\mathbb{R}^1, \mathbb{R}^2)$  be a  $k$ -graph of some solution of system (3) and let  $\Phi_t$  be the shift along the vector field  $S$ . Then the surface  $\Phi_t(\Gamma^k)$  is a  $k$ -graph of a solution of system (2).

**Example 2.** Consider the telegraph equation

$$u_{tt} - u_{xx} = au + bu_t + c, \quad (6)$$

where  $a, b, c$  are constants. This equation admits two types of dynamics:

$$\begin{cases} y_2 = \frac{y_1}{x + \alpha}, \\ z_2 = \frac{z_1}{x + \alpha} \end{cases} \quad (7)$$

and

$$\begin{cases} y_2 = \frac{2b\alpha - (x + \beta)\alpha^2}{4b^2 + 16a - \alpha^2(x + \beta)^2} \times y_1 - \frac{4\alpha}{4b^2 + 16a - \alpha^2(x + \beta)^2} \times z_1, \\ z_2 = -\frac{4a\alpha}{4b^2 + 16a - \alpha^2(x + \beta)^2} \times y_1 - \frac{2b\alpha + \alpha^2(x + \beta)}{4b^2 + 16a - \alpha^2(x + \beta)^2} \times z_1. \end{cases} \quad (8)$$

Here  $\alpha, \beta$  are arbitrary constants. The general solution of equation (7) is

$$\begin{cases} y(x) = C_3 + C_4(x + \alpha)^2, \\ z(x) = C_1 + C_2(x + \alpha)^2, \end{cases} \quad (9)$$

and the general solution of equation (8) is

$$\begin{cases} y(x) = \frac{1}{2}C_2x^2 + C_3x + C_4, \\ z(x) = \frac{1}{8\alpha} (x(C_2\beta - C_3)(2\beta + x)\alpha^2 + (8C_1 + 2bx^2C_2 + 4bC_3x)\alpha - 32 \left(a + \frac{b^2}{4}\right) C_2x). \end{cases} \quad (10)$$

Here  $C_1, \dots, C_4$  are arbitrary constants. Applying the shift transformations  $\Phi_t$  to the obtained general solutions, we obtain particular solutions of equation (6). For example, the function

$$\begin{aligned} u(t, x) = & -1 + \frac{1}{10} \left( \frac{5}{2}x^2 + 5 + (10x + 1 - t)\sqrt{5} \right) e^{-\frac{1}{2}(t\sqrt{5}-1)} + \\ & + \frac{1}{10} \left( \frac{5}{2}x^2 + 5 + (-10x - 1 + t)\sqrt{5} \right) e^{\frac{1}{2}(t\sqrt{5}-1)} \end{aligned} \quad (11)$$

is a solution of equation (6). It corresponds to solution (10) with  $a = b = c = 1$ ,  $\alpha = 1, \beta = 0$  and  $C_1 = 0, C_2 = 1, C_3 = 0, C_4 = 0, C_5 = 0$ .

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