

# Conformal mappings and invariance of the Einstein equations

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Let  $\psi$  be a spinor field in an  $n$ -dimensional space-time  $(V^{1,n-1}, g)$ . The Belinfante–Rosenfeld stress-energy tensor of the spinor field is calculated by the formula [1]

$$T_{jk} = \frac{i}{2} (\bar{\psi} \gamma_{(j} \nabla_{k)} \psi - (\nabla_{(j} \bar{\psi}) \gamma_{k)} \psi). \quad (1)$$

It should satisfy the Einstein equations [1]

$$R_{ij} - \frac{1}{2} R g_{ij} = -8\pi G T_{ij}, \quad (2)$$

where  $g_{ij}$ ,  $R_{ij} = R_{ij\alpha}^{\alpha}$  and  $R = R_{ij} g^{ij}$  are the metric, the Ricci tensor and the scalar curvature respectively of a space-time  $(V^{1,n-1}, g)$ . Obviously,  $R_{ij}^h$  is the curvature tensor of  $(V^{1,n-1}, g)$ .

Let us consider a conformal mapping  $f : (V^{1,n-1}, g) \rightarrow (\tilde{V}^{1,n-1}, \tilde{g})$ , i. e. a diffeomorphism between the manifolds whose metrics are related as

$$\tilde{g}_{ij}(x) = e^{2\varphi(x)} g_{ij}(x),$$

where  $\varphi(x)$  is a function of the  $x$ 's. The following relations hold [3, 4]

$$\begin{aligned} \tilde{R}_{ijk}^h &= R_{ijk}^h + \delta_k^h \varphi_{ij} - \delta_j^h \varphi_{ik} + g^{hl} (\varphi_{lk} g_{ij} - \varphi_{lj} g_{ik}) + (\delta_k^h g_{ij} - \delta_j^h g_{ik}) \Delta_1 \varphi, \\ \tilde{R}_{ij} &= R_{ij} + (n-2) \varphi_{ij} + (\Delta_2 \varphi + (n-2) \Delta_1 \varphi) g_{ij}, \\ \tilde{R} &= e^{-2\varphi} (R + 2(n-1) \Delta_2 \varphi + (n-1)(n-2) \Delta_1 \varphi). \end{aligned} \quad (3)$$

Here  $\varphi_i = \partial_i \varphi$ ,  $\Delta_1 \varphi = \varphi_i \varphi_j g^{ij}$ ,  $\varphi_{ij} = \nabla_j \varphi_i - \varphi_i \varphi_j$ ,  $\Delta_2 \varphi = \nabla_j \varphi_i g^{ij}$ . We know that invariance of the Dirac equations requires that a spinor field should be transformed according to the formula [1]

$$\psi' = e^{s\varphi(x)} \psi.$$

The number  $s = \frac{1-n}{2}$  is referred to as the conformal weight of a spinor field  $\psi$ .

Then the Belinfante–Rosenfeld stress-energy tensor (1) of a spinor field is transformed within a conformal mapping according to the formula

$$\tilde{T}_{jk} = e^{(2s+1)\varphi} T_{jk}. \quad (4)$$

Note that for a massless spinor field [2, 5]

$$\tilde{T} = \tilde{T}_{jk} \tilde{g}^{jk} = T = T_{jk} g^{jk} = 0.$$

It follows from (2), (3) that the function  $\varphi$  must satisfy the system (massless case)

$$\nabla_j \varphi_i = \varphi_i \varphi_j - \frac{1}{2} g_{ij} \Delta_1 \varphi - \frac{8\pi G}{n-2} (\tilde{T}_{ij} - T_{ij}),$$

or, taking into account (4),

$$\nabla_j \varphi_i = \varphi_i \varphi_j - \frac{1}{2} g_{ij} \Delta_1 \varphi - \frac{8\pi G T_{ij}}{n-2} (e^{(2s+1)\varphi} - 1). \quad (5)$$

Conditions of integrability of (5) are [6]

$$\varphi_{\alpha} R_{ijk}^{\alpha} + \frac{8\pi G}{n-2} ((\nabla_k T_{ij} - \nabla_j T_{ik}) (e^{(2s+1)\varphi} - 1) + (\varphi_k T_{ij} - \varphi_j T_{ik}) (2s+1) e^{(2s+1)\varphi}).$$

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